

On Electro-Magnetic Inertia and Thrust

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1. Introduction

When electric or magnetic systems are charged with energy, it is to be expected that the effective mass of the system increases according to Einstein's mass-energy equivalence $W = mc^2$ (we have used the symbol W to represent energy instead of the usual E to prevent confusion with the use of E as electric field), i.e. the mass increase is given by

$$m = \frac{W}{c^2} \quad (1)$$

Many texts refer to this as *electrostatic mass* or *electromagnetic mass*, which for practical purposes is negligible because of the large magnitude of c^2 .

However there are other inertial contributions induced dynamically which could be called *electrodynamical mass*. This paper deals with these dynamic situations where the mass contribution can be positive or negative. The derivation of this mass involves computation of electrodynamic forces created by acceleration from which the effective inertial mass automatically follows. However it is not necessary to think in terms of inertial mass, those forces exist in their own right and could be put to good use. The final part of this paper considers systems where those forces appear.

2. Electrically Induced Inertia.

Consider two electrically charged spheres separated by a distance r , Q_2 being fixed while Q_1 is accelerated, Figure 1.

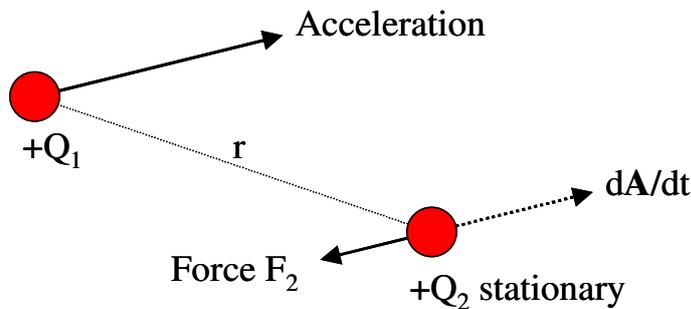


Figure 1. Force induced by acceleration

The moving charge Q_1 creates at Q_2 a vector magnetic potential field \mathbf{A} proportional to its velocity given by

$$\mathbf{A} = \frac{\mu_0 Q_1 \mathbf{v}}{4\pi r} \quad (2)$$

and since Q_1 is accelerating the \mathbf{A} field is changing with time

$$\frac{d\mathbf{A}}{dt} = \frac{\mu_0 Q_1}{4\pi r} \cdot \frac{d\mathbf{v}}{dt} \quad (3)$$

The changing \mathbf{A} field creates an electric field \mathbf{E} given by

$$\mathbf{E} = -\frac{d\mathbf{A}}{dt} \quad (4)$$

so that Q_2 endures a force $Q_2\mathbf{E}$ from that induced \mathbf{E} field

$$\mathbf{F}_2 = -\frac{\mu_0 Q_1 Q_2}{4\pi r} \cdot \frac{d\mathbf{v}}{dt} \quad \text{Newtons} \quad (5)$$

This is simply the near-field radiation from one accelerating charge affecting a second nearby charge.

If we move our reference frame from that fixed in Q_2 to one fixed in Q_1 , we find that Q_2 is accelerating with reference to Q_1 , so that Q_1 endures a force

$$\mathbf{F}_1 = +\frac{\mu_0 Q_1 Q_2}{4\pi r} \cdot \frac{d\mathbf{v}}{dt} \quad \text{Newtons} \quad (6)$$

see Figure 2.

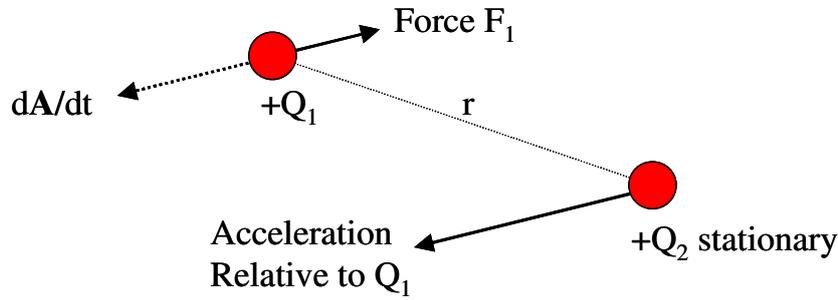


Figure 2. Induced Reaction Force

Equations (5) and (6) may be considered to be manifestations of Newton's third law of action and reaction. Note that (6) is an inertial effect occurring at Q_1 , in this case (of two like charges) a negative form of inertia with a force that aids the acceleration rather than opposing it. If Q_2 were of opposite polarity (6) would predict positive inertia, a force that opposes the acceleration, i.e. the sort of inertia we are familiar with when accelerating mass.

Equation (6) can be manipulated into a version using the electrostatic potential ϕ_E from Q_2 which is given by

$$\phi_E = \frac{Q_2}{4\pi\epsilon_0 r} \quad \text{volts} \quad (7)$$

yielding

$$\mathbf{F}_1 = +\frac{Q_1\phi_E}{c^2} \cdot \frac{d\mathbf{v}}{dt} \quad \text{Newtons} \quad (8)$$

where we have also used $\mu_0\epsilon_0 = \frac{1}{c^2}$. We could surround Q_1 with a "galaxy" of fixed charges, and still use (8) where ϕ_E is then the sum of potentials from that galaxy. The inertial effect given by (8) could be represented by an induced inertial mass m_i which adds to the actual mass of Q_1

$$m_i = -\frac{Q_1\phi_E}{c^2} \quad \text{Kilograms} \quad (9)$$

$Q_1\phi_E$ can be recognised as the *potential energy* of Q_1 , i.e. the energy to move Q_1 out of the potential ϕ_E into a far region of space where the potential is zero. In this case the negative sign indicates that Q_1 would eventually give up that energy as it is repulsed by the galaxy of positive charges. Thus in (9) we have arrived at Einstein's mass-energy equivalence (1) by the use of classical electromagnetic theory, but note this allows both positive and negative inertia to be induced.

So far we have not considered retardation effects due to the finite propagation delay. This certainly affects the force on Q_2 (or on all the other charges in our "galaxy") because that force is time-delayed by a time r/c . But the electric potential ϕ_E through which Q_1 accelerates is not part of that time delay, that potential already exist so the *inertial* effect at Q_1 (8) or (9) is instantaneous.

3. Magnetically Induced inertia

There are many duals between electric and magnetic effects where equations follow similar formats, in particular the forces between point charges and between point magnetic poles. Although magnetic poles don't really exist as such, the formulae are still useful for predicting magnetic effects. Thus by analogy to (9) it is to be expected that a magnetic pole of strength Q_m amp-meters, when placed in a scalar magnetic potential of ϕ_m amps, will inherit a magnetically induced inertia which can be represented by an inertial mass m_m kilograms given by

$$m_m = \frac{\mu_0 Q_m \phi_m}{c^2} \quad \text{Kilograms} \quad (10)$$

Similar to the electric case, $\mu_0 Q_m \phi_m$ is the potential energy of the magnetic pole Q_m , i.e. the energy involved in moving it from the potential ϕ_m to a far region of space where the potential is zero. And as in the electric case the induced inertial mass can be positive or negative.

3.1. Translatory Acceleration

Isolated magnetic poles don't exist, so we must examine a pair of magnetic poles representing a magnetic dipole, a permanent magnet. For translatory acceleration where we must consider the total inertia of the magnet, in a uniform scalar potential the induced inertia of the pair would cancel, it requires the scalar potential on one pole to be different in value from that on the other pole, viz. there must be an H field component along the dipole axis (rotation is considered separately later). Denoting this component as H_μ , we can deduce the induced inertial mass of the system to be represented by

$$m_m = \frac{\mu_0 Q_m H_\mu l}{c^2} \quad \text{Kilograms} \quad (11)$$

where l is the length of the dipole. Note that this is an *inertial* effect separate from the usual quasi-static magnetic forces on the poles that come from the gradient of the potential (i.e. an H field) producing an angular alignment force on the magnet, or the spatial gradient of the H field producing a translatory force on a stationary magnet. Here we are dealing with forces produced when the magnet is accelerated.

In the case of a permanent magnet the pole strength Q_m is given by

$$Q_m = Ms \quad \text{Amp-metres} \quad (12)$$

where s is the surface area of the pole face and M is the magnetization. Hence (11) becomes

$$m_m = \frac{\mu_0 M H_\mu V}{c^2} \quad \text{Kilograms} \quad (13)$$

where V is the volume of the magnet. Now since

$$\mu_0 M = B_{sat} \quad (14)$$

where B_{sat} is the saturation flux density we can rewrite (13) as

$$m_m = \frac{B_{sat} H_\mu V}{c^2} \quad \text{Kilograms} \quad (15)$$

Taking a typical small magnet volume as 10 cm^3 (10^{-5} m^3), using 1 Tesla magnets close to each other so that $H \approx 1/\mu_0$ we can estimate the mutually induced inertias as $8.83\text{E-}17 \text{ Kg}$. That is so small as to be negligible.

3.2. Constant Rotation

For rotation about the dipole centre, because the induced inertial mass (10) at opposite poles is of opposite polarity, even in a uniform potential we get an unbalance with respect to centrifugal forces. Again the effect is found to be small except for the special case detailed next.

Although the magnetic **field** from the Earth is quite weak, the same cannot be said for the magnetic scalar **potential** ϕ_m . The magnetic Earth can be modelled as a magnetic dipole at its centre whose dipole moment μ_E is $8.24 \times 10^{22} \text{ amp-m}^2$. The scalar magnetic potential for a dipole is given by

$$\phi_m = \frac{\mu_E \cos \theta}{4\pi r^2} \text{ amps} \quad (16)$$

where r is the radial distance and θ is measured from the dipole axis. The radius of the Earth is 6.37×10^6 meters. At the equator where θ is 90° the potential is zero, but at a latitude of say 30° where $\cos(\theta)$ is 0.5 we get a value of 8×10^7 amps, and at the Earth's magnetic poles the value would be 1.6×10^8 amps. The presence of this huge scalar potential is not taught hence is little known, perhaps because until now there has been no method of using it to practical purpose. For calculation purposes the notional scalar potential will be taken as that 8×10^7 amps value.

Putting (12) and (14) into (10) we get for the inertial mass induced at the poles of a permanent magnet by the presence of ϕ_m

$$m_m = \pm \frac{B_{sat} s \phi_m}{c^2} \quad (17)$$

Note that $B_{sat} s \phi_m$ is the potential energy of each magnet pole taken in isolation. Thus for a 1 Tesla magnet with $s=4\text{cm}^2$ pole faces the potential energy for each pole has a magnitude of 32 KJ in the notional Earth's scalar potential. The induced mass appears as an increase in the real mass at one pole and a decrease at the other pole. If we rotate the magnet about its centre at constant speed, we obtain a dynamic unbalance where the radial centrifugal forces do not cancel, there is a net force in one direction. From the centrifugal force equation

$$F = m r \omega^2 \quad (18)$$

we find that the imbalance creates an anomalous force

$$F_A = \frac{B_{sat} s l \phi_m \omega^2}{c^2} \quad \text{Newtons} \quad (19).$$

The radial force vector F_A is of course rotating with the magnet and points along its axis. Converting the product of pole area s and length l into the volume V , using (14) for B_{sat} and since the product MV is the dipole moment of the magnet we get the satisfying result that applies to any dipole of moment μ

$$F_A = \frac{\mu_0 \mu \phi_m \omega^2}{c^2} \quad \text{Newtons} \quad (20)$$

Taking a 1 Tesla magnet of pole area $s=4\text{cm}^2$ and of length 20cm, if we rotate this about its centre at say 3000rpm then from (19), in the notional Earth's magnetic potential, we find an induced unbalancing force of 7×10^{-9} Newtons. Again this is rather small to be of any practical use.

However it is well known that high-speed rotation of a magnetic dipole can be obtained using high permeable material excited by coils. Thus a sphere or disc of permeable within a pair of coils at right angles and driven in phase quadrature, can create rotation speeds of its magnetization of MHz and higher. If the above magnet were simulated in this manner at a rotation speed of 1MHz, the induced force anomaly would be in the order of 2.8 Newton, which is much more respectable. Get rotation at 10MHz and we have a 28Kg force!!

The force unbalance will create the equivalent of a rotating force vector acting on the whole body. Thus along a given axis this will constitute a vibratory force at the rotation frequency, *this force being induced because of the presence of the earth's magnetism in the form of its local scalar potential*. We have the technology for measuring such a vibratory force and this experiment would seem to be worth doing in order to validate the theory.

3.3. Angular Acceleration

The dynamic unbalance of mass at the poles of the magnet also creates an anomalous force on the system when it is accelerated angularly. We find that in this case the anomalous force F_A is given by

$$F_A = \frac{B_{sat} s l \phi_m}{c^2} \cdot \frac{d\omega}{dt} \quad \text{Newtons} \quad (21)$$

which can be expressed for any dipole μ as

$$F_A = \frac{\mu_0 \mu \phi_m}{c^2} \cdot \frac{d\omega}{dt} \quad \text{Newtons} \quad (22)$$

In this case the radial force acts at right angles to the dipole axis. This could account for anomalous translatory impulse forces when a magnet suddenly receives a rotation impulse, perhaps it explains the Steorn effect. The anomaly would disappear at the equator, and be a maximum at the Earth's magnetic poles.

3.4. Potential Thrust Motor

A series of alternate polarity rotation impulses (i.e. alternate angular acceleration and deceleration) timed to occur each 180° of rotation could create a net unidirectional

force. Again at speeds associated with moving masses this would be quite small, but if engineered to occur within a disc of magnetically permeable material using coils driven by appropriate waveforms the force could be significant. Could this form the basis for a thrust motor using the Earth's scalar magnetic potential? It would seem to be something worth exploring.

4. Systems using acceleration forces

As stated above a magnetic dipole can be created within a sphere or disc of magnetically permeable material that is within an energized coil. Having two coils at right angles and fed with AC at 90° phase difference will create a rotating dipole at angular speeds far in excess of those achieved mechanically. Any advantage to be gained from the electrodynamic forces is damped by the real mass of the sphere or disc. However there is another possibility for creating such a high-speed rotation where the real mass is much lower. A ring core having two toroidal windings each occupying only half the core can create N and S poles that are diametrically opposite as shown in figure 3.

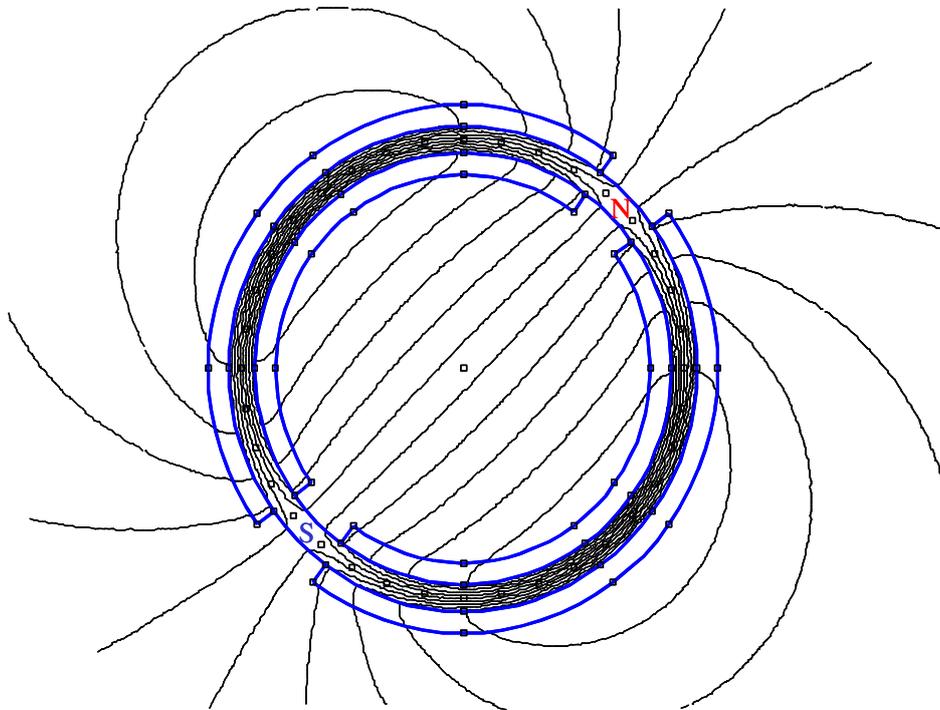


Figure 3. FEMM Plot of ring core with two bucking windings

The real mass can be kept low by having a core of minimum cross section, e.g. the core is now simply a loop of iron wire. Application of two more coils at 90° to the first pair and fed with a 90° phase difference will induce a rotation of that dipole at the drive frequency, i.e. the poles move along the wire. At sufficient rotation speeds the inertial effect in the Earth's scalar magnetic potential could invoke radial centrifugal-like forces on the wire.

When dealing with the forces on those N and S poles we have to consider separately the two cases (a) where the force is at right angle to the wire and the wire tries to move and (b) where the force is along the wire and the pole can move within the wire. In the first case the actual movement will be restricted by the real mass and the

support structure while in the second case there is no such restriction. It is also instructive to note the difference between positive and negative inertia. Positive inertia where the force opposes the acceleration we are all familiar with since mass behaves in this way. Negative inertia where the force aids the acceleration is clearly an unstable state where the thing being accelerated (in our case a magnetic pole) continues to accelerate at ever increasing rate and, given no restraining influence, would disappear into infinity!

4.1. Pole acceleration along a Fe wire

Let us imagine a Fe wire hoop that has induced in it a S pole by the presence of a permanent magnet as shown in figure 4.

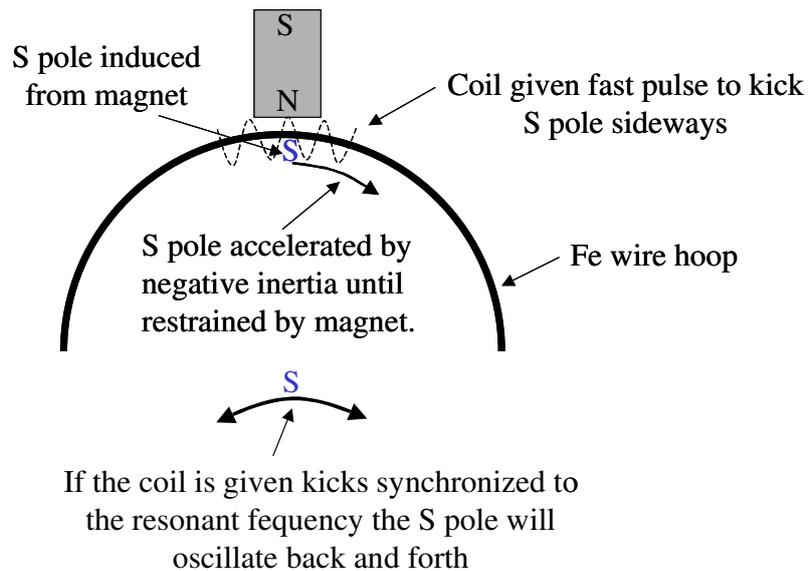


Figure 4. Movement of unstable pole on Fe Hoop

If we now try to move the pole by supplying a small coil around the Fe with a fast pulse, in the north hemisphere of the Earth the S pole will endure a force supporting the acceleration, it is unstable. Hence that pole will quickly accelerate away from the magnet, only to be stopped by the attractive force between the pole and the magnet. That attractive force will then accelerate the pole back in the opposite direction, again aided by the negative inertia. By a series of “kicks” at the resonant frequency that unstable pole can be made to oscillate back and forth perhaps by a significant amount.

4.2. Pole acceleration transverse to a Fe wire

The wire hoop is circular, so the oscillatory movement along a circular path will create centrifugal radially inward acceleration, so now we have the case where an inertial force is invoked transverse to the wire. This will cause the wire to move and since the translation is oscillatory so will be that movement. Now we have a wire vibrating within a magnetic field and this will induce voltage into the wire. The wire could be a multi turn loop of thin Fe and the voltage induced could be used to drive current through a load. Normally that current would load the oscillatory driving force and any power drawn would come from that mechanical drive, but here we have that

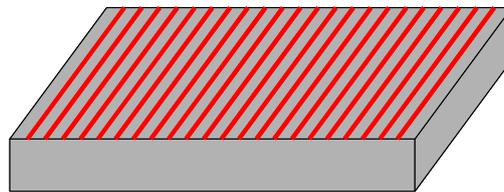
mechanical drive derived from the Earth's scalar magnetic potential, initiated by a series of kicks.

We could create two unstable poles by having another magnet at a diametrically opposite position. That arrangement looks very similar to Steven Mark's TPU.

4.3. Surface pole on magnet slab

By having a magnetizing coil closely wound onto a ferrite magnet with appropriate wire spacing it is possible to obtain a striped pattern of magnet poles. Having such a conditioned magnet within a solenoidal coil we can create sideways kicks to those poles, and the array of unstable poles can be made to oscillate in a similar manner to that described above, figure 5.

Magnet conditioned to have striped pattern of poles



Unstable pole array made to oscillate

Figure 5. Oscillatory array of surface poles

That oscillatory movement can induce voltage into a coil closely wound onto the magnet with wire spacing matched to the magnetic stripes. Could this explain the Floyd Sweet VTA?

5. Conclusion.

It has been shown that there is the possibility of invoking inertial forces on magnetic poles that are within a scalar magnetic potential. The Earth's scalar potential can be as high as 1.6×10^8 amps which should allow significant forces to be observed. The use of this feature for obtaining thrust has been examined. The possibility of obtaining anomalous energy has also been looked at and it seems likely that this inertial feature could explain the Steven Mark's TPU and the Floyd Sweet VTA.